# Effect of scatter coincidences, partial volume, positron range and non-colinearity on the quantification of FDOPA Patlak analysis

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## ABSTRACT

**Introduction:** The key characteristics of positron emission tomography (PET) are its quantitative capability and its sensitivity, which allow the in vivo imaging of biochemical interactions with small amounts of tracer concentrations. Therefore, accurate quantification is important. However, it can be sensitive to several physical factors. The aim of this investigation is the assessment of the effect of physical effects, such as: scatter coincidences, partial volume, positron range and non-colinearity on the quantification of FDOPA uptake using PET.

**Methods:** The SimSET Monte Carlo package was employed to simulate acquisitions of the PET/CT Siemens Biograph scanner. The study was performed with a numerical brain model obtained from the CT scan of a commercial striatal phantom. Theoretical pharmacokinetic values were simulated. The simulations were carried out with and without scatter, positron range and non-colinearity effects. The OSEM algorithm from STIR library was used to reconstruct the PET data. Different correction strategies were employed in order to evaluate the effects caused by the different type of degradation on results obtained with Patlak analysis.

**Results:** The FDOPA uptake of Patlak plot increased from 70.4% of the theoretical value to 80.4%, if scatter was perfectly corrected, and it increased to 99% of the theoretical value when the partial volume correction was employed, as well. No significant improvement was found for positron range and non-colinearity effects when the partial volume correction was employed.

**Conclusions:** The results show that the compensation for scatter and partial volume degradations increases accuracy in the uptake calculation.

Key words: PET, Reconstruction, Kinetic analysis, Monte Carlo simulation.

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## **INTRODUCTION**

Positron Emission Tomography (PET) facilitates the diagnosis of a disease state at an early point and the monitoring of its evolution, vielding important biological insights, informing the development and evaluation of therapies. This is usually achieved with dynamic PET imaging which involves a timesequence of acquisitions called frames. Each frame is independently reconstructed into separate images. Compartmental models are then applied to the time radioactivity images of either anatomical regions or individual voxels in order to estimate physiological parameters. This information contains the kinetic constants of the biochemical process, which relate the total temporal concentration of the radiotracer and the input function. The kinetic constants are usually estimated by fitting the model to the time-course of the images.

For example, FDOPA (6-[<sup>18</sup>F]Fluoro-L-dopa) and <sup>11</sup>C-raclopride studies in dynamic PET are used to study the presynaptic dopamine function and reveal the status of the dopaminergic nigrostriatal pathway both at presynaptic and postsynaptic level [1]. In particular, FDOPA is an exogenous substrate of dopa-decarboxylase that is the immediate precursor of dopamine [2, 3]. The FDOPA uptake has also been found to decrease in patients with Parkinson's disease. Thus, FDOPA studies are mainly used in the examination of patients affected by Parkinson's disease [4]. The use of kinetic models instead of a static analysis provides a better objective marker of Parkinsonism or other relevant diseases [5]. The kinetic information in the FDOPA studies can be measured with a two-tissue three-compartments model, so that the total measured radioactivity can be subdivided into reversible compartments (i.e. tracer molecules are freely exchanged between tissue and plasma) and irreversible (i.e. tracer molecules are trapped in these compartments). If the tissue with reversible kinetic behaviour of the tracer is in steady state, the overall uptake constant of FDOPA from plasma to the irreversible compartment can be estimated as the slope of a linear regression. This linear model is called Patlak plot [6].

One problem associated with quantitative PET analysis is that accurate quantification depends on several factors. There are many technical and physical factors to be considered in this analysis, which result in inexact data. Several papers have investigated the effect of these physical factors on PET FDG studies. In regard to the coincidence detection process random coincidences can occur when two or more photons from different annihilation events are detected within the same time window. This effect is addressed by estimating the random coincidences and then considering them in the reconstruction. The rate of random coincidences is proportional to the singles rate and the coincidence time window [7]. On the other hand, there are effects related to the photon interactions in human tissue such as the attenuation effect. This effect is due to the loss of true coincidence events (photon absorption or scattering) and is addressed through the use of density maps from the current PET/CT scanners [8]. Some of these physical effects have been also evaluated in dynamic FDG PET studies, such as the reconstruction method [9, 10] and scatter and randoms contribution [11, 12]. It is also essential to note several temporal frames are normally necessary for kinetic analysis, but most of the frames have low number of counts and consequently poor signal to noise ratio [Author et al 2011, 13]. In this context, another paper has evaluated several parameter estimation methods (including the Patlak plot) that are necessary for the determination of the cerebral metabolic rate of individual kinetic rate constant parameters [14]. The correction for other physical effects such as positron-range and non-colinearity is not part of the reconstruction process and consequently may affect the accuracy of the clinical studies. An attempt to inherently include such effects is the recent emerging resolution modelling (PSF) method, which is currently an active research area. Effects such as the partial volume are related to the limited spatial resolution of the imaging system [15]. In general, these effects cause erroneous estimation of the radioactivity in each volume of interest and potentially affect the estimation of the pharmacokinetic parameters [16, 17]. Most of literature has investigated the effect of different reconstruction strategies and physical factors on static FDG analyses, and only few on dynamic FDG. Furthermore, the kinetic parameters may be affected in a different way than the static quantification of linear models. Moreover, to the best of our knowledge, there are no papers aimed that investigate the influence of reconstruction strategies and physical effects in the kinetic analysis of dynamic FDOPA PET studies.

The aim of this work is the assessment of the physical effects of scattered coincidences [18], partial volume effect [19], positron range and non-colinearity [20, 21] on the quantification of striatum and reversible compartment in FDOPA studies and consequently their impact on the physiological parameters estimated by the Patlak plot. Once the influence of each effect is classified, our work aims to test a new procedure for carrying out more reliable quantification studies.

The assessment is carried with Monte Carlo simulation, which is an essential tool for understanding and evaluating these physical effects as it allows examine complex systems than we otherwise can. Several Monte Carlo packages have been developed and are widely used. One recent

example was the comparison of the performance of several commercial scanners [22]. In our investigation we used the SimSET Monte Carlo package [23].

## Theory

## Compartment model

A two-tissue three-compartments model for FDOPA is described and the theoretical concentration variations along time within each brain region are obtained analytically in order to simulate a dynamic FDOPA study. The mean radioactivity concentrations within the striatum and the other regions are obtained from reconstructed images and the overall uptake constant of FDOPA is estimated by linear regression. In the two-tissue three-compartments model [24] shown in Figure 1, it is assumed that the system can be subdivided into reversible and irreversible chemical compartments. The input function is assumed to be the arterial plasma time activity curve (TAC), which is usually measured during a human brain study. If  $C_T^{(Irr)}(t)$  and  $C_T^{(Rev)}(t)$  are the concentrations of the labelled compounds in the irreversible and reversible compartments,  $C_{T}(t)$  is the total radioactivity in the tissue and  $C_P(t)$  is the plasma time radioactivity curve [25], the corresponding radioactivity of each compartment will be:

$$C_T^{(hr)} = \frac{K_1 k_3}{(k_2 + k_3)} * C_P$$
(Eq. 1)

$$C_T^{(\text{Re}\nu)} = \frac{K_1 k_2}{k_2 + k_3} e^{-(k_2 + k_3)t} * C_P$$
 (Eq. 2)

The symbol \* denotes the mathematical convolution operation.



Fig 1. Two-tissue three-compartments model.

## The Patlak Model

The Patlak approach can describe the behaviour of the FDOPA model when the free radiotracer in tissue has reached a steady state. This approach is valid only at sufficiently late time points, for example 50 min following the bolus passage of the tracer. Thus, the original model is simplified and  $C_T(t)$  can be

obtained in good approximation by the following linear equation:

$$C_{T} = \frac{K_{1}k_{3}}{(k_{2} + k_{3})} \int_{0}^{t} C_{P}(\tau) d\tau + \frac{K_{1}k_{2}}{(k_{2} + k_{3})^{2}} C_{P}(t) \quad (\text{Eq. 3})$$

Which after the transformations:  $b = \frac{K_1 k_2}{k_2}$ 

 $K_1k_3$ 

 $(k_{2} + k_{3})$ 

$$(k_2 + k_3)^2$$
 takes the simple form:

$$\frac{C_T}{C_P(t)} = a \frac{\Box C_P(\tau) d\tau}{C_P(t)} + b$$
 (Eq. 4)

This is known as the Patlak Plot which allows the estimation of the parameters a and b using linear regression. The parameter *a* represents the overall uptake constant of FDOPA from the plasma to the striatum and it is the main parameter of interest. The parameter b represents the free FDOPA fraction in tissue and blood.

#### Partial volume correction method

The size of the volumes of interest, if particularly small as the caudate and the putamen, has a direct impact on the measurement of the radioactivity concentration. Thus, correction for PVE is very useful in order to obtain accurate quantification of small regions as the striatum [19].

An approach to correcting the PVE in neuroreceptor imaging has been described by different groups [26, 27]. We modified the Fleming method to obtain the total striatum uptake. The methodology is based on the measurement of the total radioactivity in the striatum and the calculation of the exact volume of interest for each region. The total radioactivity in the striatum was calculated by using expanded regions of interest (ROIs) which are large enough to include the entire radioactivity that has spread outside the physical volume of the structures due to the PVE. The extension of the expanded ROIs is related to the spatial resolution of the PET system.

Thus, the expanded ROIs were obtained from the original ROIs after convolution with a Gaussian function of full-width at half maximum (FWHM) equal to the crystal size of the scanner (4 mm).

The original ROIs of the striata and other regions were defined by using a CT image of the anthropomorphic striatal brain phantom. Thus, the exact volumes of the ROIs were calculated. In

clinical studies, the striata regions could be obtained by segmenting MR images and then registered them with PET images. To quantify the studies, all reconstructed images were resampled to the highresolution space that the ROIs were defined.

Compensation for the PVE was carried out by expansion of the original striatum and then calculation of the mean radioactivity concentration within this expanded region. Then, the radioactivity concentration inside the original striatum ( $C_{Str}$ ) was calculated by subtracting the non-striatum region (i.e. reversible component) from the total radioactivity in the expanded region:

$$\mathbf{C}_{\text{Str}} = \left(\mathbf{C}_{\text{Str}}^{m} V_{\text{Str}'} - \mathbf{C}_{n\text{Str}}^{m} \left(V_{\text{Str}'} - V_{\text{Str}}\right)\right) \cdot V_{\text{Str}} \quad (\text{Eq. 5})$$

 $C_{\text{Str}}^{m}$  and  $C_{n\text{Str}}^{m}$  are respectively, the mean radioactivity concentrations within the expanded striatum and the non-striatum regions;  $V_{\text{Str}}$  and  $V_{\text{Str}}$  are the volumes of the original and the expanded striatum, respectively.

The concentrations of the labelled compounds in the irreversible  $(C_T^{(Rev)})$  and the reversible compartments  $(C_T^{(Rev)})$  were obtained from these concentrations  $(C_T^{(Rev)}(t) = C_{nStr}(t)$  and  $C_T^{(drr)}(t) = C_{Str}(t) - C_{nStr}(t)$ ) and the Patlak plot was then used to calculate the parameters of interest: *a*, *b*.

## **METHODS**

## Numerical phantom

A numerical phantom was obtained from a CT scan of the anthropomorphic striatal brain phantom (Radiological Support Devices Inc., Long Beach, CA, USA). The CT image consisted of  $256 \times 256 \times 196$ voxels with size of  $0.89 \times 0.89 \text{ mm}^3$ . Brain tissue and bone were automatically segmented by thresholding the CT image while the striatum was manually defined.

The non-uniform attenuation map was obtained by setting the appropriate attenuation coefficients to brain and bone for 511 keV (i.e. 0.096  $cm^2/g$  for water and 0.197  $cm^2/g$  for bone).

Radioactivity distribution maps were separated into two different components: the striatum and the nonstriatum (also called background) radioactivity distribution maps. Figure 2 shows a central slice of the background radioactivity distribution (left-hand side), the striatum radioactivity distribution (middle) and the non-uniform attenuation map (right-hand side).



Fig 2. Striatal phantom: non-striatum background radioactivity distribution (left), striatum radioactivity distribution (centre) and the non-uniform attenuation map (right).

#### Monte Carlo simulation

The SimSET Monte Carlo package [23] was employed to simulate fully 3D acquisitions of the PET/CT Siemens Biograph scanner (the source code library can be downloaded from: http://depts.washington.edu/simset/html/simset main. html). As a consequence, 1024 direct and indirect sinograms with 288 transaxial bins (bin size of 2.2 mm) and 288 angular positions were obtained. Simulations were carried out with or without positron range and non-colinearity effects. True and scatter coincidences were stored separately so that an ideal scatter correction (ISC) could be performed. Sinograms of non-striatum background and striatum radioactivity were generated separately. These simulations were carried out in such a way that noise was 100 times less than what is normally obtained from a real acquisition. Thus, these simulations can be regarded as approximately noise-less, providing the advantage to assess the physical effects without introducing large statistical errors.

The protocol consists of scans over a 90 min period divided into 26 frames: 1 of 40 s, 4 of 60 s, 3 of 120 s, 3 of 180 s and 15 of 300 s, as it is a standard optimized FDOPA protocol [28]. The simulation was carried out for the last 8 frames that are used in the Patlak model. Both the non-striatum and striatum simulations were multiplied by the appropriate factor and then Poisson noise was added in order to mimic the number of counts collected per frame. Then, striatum sinograms were subsequently added to the background sinograms. Once the simulation process was completed, three different sets for each of the last eight time frames were generated containing: total coincidences, trues coincidences (ISC) and trues coincidences without positron-range and noncolinearity effects (ISC + IPRC + INCC).

#### Reconstruction

The fully 3D OSEM iterative algorithm implemented in STIR library [29, 30] was used to reconstruct the 3D PET data on a reconstruction grid of 128×128×63 (pixel size of 1.7 mm and slice thickness of 2.5 mm). (The source code of STIR library can be downloaded from the webpage: http://stir.sourceforge.net). Image counts per voxel were then converted by proper scaling into kBq/ml. In order to ensure convergence, all images employed in the Patlak plot analysis were reconstructed using adequate number of iterations. This optimal number of iterations is usually object-dependent and convergence does not necessarily occur at the same iteration. The appropriate number of sub-iterations was determined by comparing the  $C_{Str}$  and  $C_{nStr}$  values obtained at consecutive iterations as shown in Eq. 6:

$$d = \frac{C_{\text{Str}} - C_{n\text{Str}}}{C_{n\text{Str}}}$$
(Eq. 6)

For the evaluation of the effect of physical degradations we employed the following correction strategies:

- *OSEM3D:* reconstruction of sinograms containing total coincidences (i.e. uncorrected data).
- *OSEM3D+ISC:* reconstruction of sinograms containing unscattered coincidences, equivalent to an ideal scatter correction (ISC).
- *OSEM3D+ISC+PVEC:* reconstruction of sinograms containing unscattered coincidences and including a partial volume effect correction (PVEC).

OSEM3D+ISC+IPRC+INCC+PVEC:

reconstruction of sinograms containing unscattered coincidences, including PVEC and without positron range and non-colinearity effects. The same PVEC method was used with and without PR and NC (i.e. ideal PR correction, IPRC and ideal NC correction, INCC) in order to evaluate how the PVEC affects the quantitative values.

#### RESULTS

## Time activity curves (TACs)

Firstly, the theoretical TACs were obtained to assign the radioactivity distribution of the numerical phantom used for the simulation. These theoretical TACs describe the FDOPA concentration variation versus time within each brain region and they are derived analytically using kinetic constants from literature ( $K_1 = 0.0403 \text{ ml/min/ml}, k_2 = 0.0342 \text{ min}^{-1}$ and  $k_3 = 0.0124 \text{ min}^{-1}$  [31] that correspond to the Patlak slope  $a = 4.965 \times 10^{-8} \text{ min}^{-1}$  and intercept b =0.635). The input function was analytically derived as the sum of three exponentials, as proposed by Yu *et al* [32]. The theoretical TACs, displayed in Figure 3, represent the background and striatum radioactivity distribution of the numerical phantom which is used for the simulation. The radioactivity concentration in striatum is obtained by adding the concentrations of the reversible and irreversible compartments ( $C_{Str}(t)$ ) =  $C_T^{(Rev)}(t) + C_T^{(Irr)}(t)$ , while the radio-activity concentration in non-striatum corresponds to the reversible compartment (  $C_{nStr}(t) = C_T^{(Rev)}(t)$  ). The reversible compartment is in steady state of the final eight frames, which correspond to the last 40 min of the scanning and approximately 50 min after the administration. The radiotracer background radioactivity in the blood volume (i.e. approximately 3-5%) was not considered in this simulation setup. This approximation is not applicable to the radioactivity in the plasma as it is part of the kinetic analysis.



Fig 3. TACs (non-striatum and striatum activities) and the input function measured from the human brain study.

#### *Number of iterations*

Once carried out the simulations, the number of iterations for the reconstruction was calculated in order to ensure convergence. Figure 4 shows  $\delta C_{Str}/C_{nStr}$  as a function of the number of iteration. These values were obtained from the reconstructed images by means of OSEM3D (4 subsets) and the results show that after several iterations convergence was achieved, which is important for accurate quantification [Author et al, 2008]. For this reason, all images employed in the Patlak plot analysis [33] were reconstructed using 60 iterations and 4 subsets.

#### Reconstruction and quantification

TACs corresponding to the reversible compartment and the striatum were derived from the reconstructed images by using the different reconstruction strategies. Figure 5 shows the comparison between the theoretical and the calculated striatum TACs for OSEM3D (uncorrected data), OSEM3D + ISC (unscattered), OSEM3D + ISC + PVEC (unscattered



**Fig 4.** Relative difference between  $C_{Str}$  and  $C_{nStr}$  obtained with OSEM (4 subsets) at consecutive sub-iterations.



**Fig 5.** Theoretical TACs for striatum (light green curve) and reversible region (dark blue curve) and calculated FDOPA striatum uptake for uncorrected; unscattered; unscattered with PR and NC effects and PVEC; and unscattered, but without PR and NC effects, including PVEC.

with PVEC) and OSEM3D + ISC + IPRC + INCC + PVEC (i.e. unscattered without PR and NC effects and with PVEC). No differences between the correction strategies in the reversible region ( $C_{nStr}$ ) were found and the calculated radioactivity levels were close to the theoretical values. This can be explained by the fact that the reversible region is much larger than the irreversible region.

Thus, even if some radioactivity concentration has migrated to another voxels due to partial volume effect, positron range, non-colinearity and scatter, the overall activity of the large reversible region is not affected. Nevertheless, both OSEM3D and OSEM3D + ISC clearly yield under-estimations of the radioactivity ( $C_{Str}$ ) in striatum TACs. PVEC recovered the value when the PR and NC effects were included in the simulation but there was about 4% bias observed when these effects were not included in the simulation. This result indicates that the specific FWHM used for the PVEC accounted also for PR and NC effects.

Figure 6 shows the relative differences of the radioactivity level between the calculated TACs and the theoretical TACs for the striatum region. This graph points out that the reconstruction from the uncorrected data provided values that underestimated the theoretical radio-activity levels by about 12% to 21%, whereas the reconstructed images from the unscattered data were underestimated by 8% to 18%.



**Fig 6.** Relative difference of the radioactivity level between the calculated and the theoretical TACs for the striatum.

#### Patlak analysis

Table 1 shows the FDOPA kinetic parameters obtained by means of different correction strategies and the results are compared with the theoretical values. It is shown that the slope (i.e. uptake constant of FDOPA) is clearly underestimated when ISC and PVEC is not considered (bias: 30% for OSEM3D, and 19% for OSEM3D+ISC). The results demonstrate also that both OSEM3D+ISC+IPRC+INCC+PVEC yield accurate estimations of the uptake parameter (0.8% and 4%).

#### Table 1. FDOPA kinetic parameters

Irreversible Compartment	Slope (min <sup>-1</sup> )	Intercept
Theoretical	0.107	0.634
OSEM3D+ISC+IPRC+INCC +PVEC	0.112	0.544
OSEM3D+ISC+PVEC	0.106	0.600
OSEM3D+ISC	0.0864	0.666
OSEM3D	0.0756	0.733
OSEM3D	0.0756	0.733

#### **DISCUSSION**

The effect of scatter is due to the fact that the scatter fraction is relatively high in 3D PET mode acquisition. The results show 4% higher bias when we did not include the PR and NC effects. This finding was not originally expected, and it can be explained by the fact that the size of the expanded ROI, which determines the success of the PVEC method, depends on the FWHM of the Gaussian function used at the convolution step. This value is related to the spatial resolution of the scanner but it is not straightforward to know which the most appropriate value for each case is and this is the main limitation of the specific PVEC approach. In the particular experiment we used a FWHM equal to the crystal detector size, however it might have been more accurate to obtain the value of the FWHM from experimental PSF measurements. Nevertheless, the observed bias is not as large as without performing the PVEC at all.

Table 1 shows that when OSEM3D + ISC + PVEC is used PR and NC effects are also compensated. The results show that the PVEC method based on expanded regions of interest can be used for FDOPA quantification, however it must be noted that the specific method assumes uniform radioactivity background. Thus, if reconstruction included modelling of the PSF, it might have improved more the quantification performance, in particular when small brain structures are examined [34].

Finally, one limitation of our study is that it was applied only in rather simple simulated data. Future investigations will focus on real data and more realistic simulations that will include effects such as motion [35], more complicated kinetic models as for example the spectral analysis or compartmental models. Moreover, synergistic information from MR or CT data for the correction of these phenomena is another potential direction to minimize their effects, improve quantitative accuracy and increase resolution [36].

#### CONCLUSION

The performance of various physical correction approaches was investigated. The methodology was based on Monte Carlo simulations allowed the assessment of the impact of different physical effects on the FDOPA Patlak analysis.

Our work suggests careful care in dealing with the quantitative values obtained from kinetic PET studies because these values can be substantially affected by different physical effects such as scatter, positron range, non-colinearity and partial volume effect. The results indicate that correction for image degrading factors is essential for accurate quantitative analysis of dynamic FDOPA PET studies. In this regard, a simple partial volume correction method is reported, which can be easily adopted by nuclear physicians in clinical routine. We recommend the use of this method in practice of nuclear medicine.

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